

VIOLATION OF BELL'S INEQUALITIES WITH A LOCAL THEORY OF PHOTONS

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Received February 26, 1996; revised July 9, 1996

We use a local theory of photons purely as particles to model the single-photon experiment proposed by Tan, Walls, and Collett. Like Tan *et al.* we are able to derive a violation of Bell's inequalities, but our local probabilistic theory does not use any specific quantum mechanical assumptions or calculations.

Key words: Bell's-inequalities, EPR, photons, hidden-variables, stochastic-models.

1. INTRODUCTION

The present paper is part of a research program which concerns a foundational analysis of phenomena usually described by quantum electrodynamics (QED) [6-8,9]. Our previous papers give a particle theory for diffraction of light and the Casimir effect. The present paper is focused on another foundational topic. It remains to be seen how far the program we have undertaken can be carried.

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A probabilistic theory of photons with well-defined trajectories is assumed. The wave properties come from the expectation density of the photons. The photons are also regarded as virtual, because they are not directly observable, including their annihilation of each other (see assumptions below). What can be detected is the interaction with matter. The meaning of *virtual* used here is not the same as in QED. In summary, our basic assumptions are:

- Photons are emitted by harmonically oscillating sources.
- They have definite trajectories.
- They have a probability of being scattered by matter.
- Absorbers, like sources, are periodic.
- Photons have positive and negative states (+-photons and --photons) which locally interfere, when being absorbed.

The expected density of \pm -photons emitted at t in the interval dt is given by

$$s_{\pm}(t) = \frac{A_s}{2}(1 \pm \cos \omega t), \quad (1)$$

where ω is the frequency of a harmonically oscillating source, A_s is a constant determined by the source, and t is time. We used $\frac{1}{2} \pm \frac{1}{2} \cos(\omega t)$ rather than $\cos(\omega t)$, to have a density that is nonnegative for all t and is between 0 and 1. If a photon is emitted at t' , $0 \leq t' \leq t$, then at time t the photon has traveled (with speed c) a distance r , where

$$t - t' = \frac{r}{c}. \quad (2)$$

The conditional spacetime expectation density of \pm -photons for a spherically symmetric source with given periodicity ω is

$$h_{\pm} = \frac{A}{8\pi r^2} \left(1 \pm \cos \omega \left(t - \frac{r}{c}\right)\right), \quad (3)$$

where A is a real constant.

The scalar field defined in terms of the expectation density $h_{\pm}(t, r|\omega)$ is

$$\mathcal{E} = \varepsilon_0 \frac{h_+ - h_-}{\sqrt{h_+ + h_-}}, \quad (4)$$

where \mathcal{E}_0 is a scalar physical constant. Using (3), (4) may be rewritten for a spherically symmetric source as

$$\mathcal{E} = \mathcal{E}_0 \sqrt{\frac{A}{4\pi r^2}} \cos \omega \left(t - \frac{r}{c} \right). \quad (5)$$

Applying the standard definition of average intensity, we get the expected result

$$I = \langle \mathcal{E}^2 \rangle = \frac{\mathcal{E}_0^2 A}{8\pi r^2}. \quad (6)$$

Note that the standard bracket notation is used for time averaging, i.e., taking an expectation with respect to t .

Since the absorber, or photodetector, behaves periodically with a frequency ω , the probability p_X of absorbing a photon in detector X is given by

$$p_X = \frac{C}{2} (1 + \cos(\omega t + \psi)), \quad (7)$$

where ψ is an arbitrary phase that can be randomized.

The expected number $E_t(X\pm)$ of each type of photon absorbed by detector X is the time-averaged product

$$E_t(X\pm) = \langle h_{\pm}^X(\theta) p_X(\psi) \rangle, \quad (8)$$

where $h_{\pm}^X(\theta)$ is the expected density of photons (with a phase θ), and ψ is a phase on detector X . Note that $E_t(X+)$, for example, is a random variable that is a function of θ and ψ . When we take expectation with respect to the distribution of θ we use subscripts to make clear that the expectation is with respect to θ . The averaging is required because an absorption of an individual photon by an atom of a photodetector takes on average several orders of magnitude longer than the mean optical period of the photons, both theoretically and experimentally [4].

As we previously assumed, during the process of absorption, photons with different states (positive and negative) annihilate each other. So, the expected number of photons to be detected in each detector X is:

$$E_t(X) = |E_t(X+) - E_t(X-)|. \quad (9)$$

We present here a violation of Bell's inequalities [1] [2] [3] with a local description of photons.

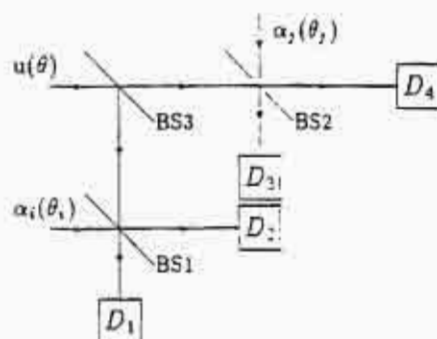


Fig. 1. Proposed experimental configuration.

2. EXPERIMENTAL CONFIGURATION

We are interested in the experimental setup proposed in [10] and also discussed in [11]. The scheme uses two coherent sources $\alpha_1(\theta_1)$, with phase θ_1 , and $\alpha_2(\theta_2)$, with phase θ_2 , and a third source to be studied, $u(\theta)$, with unknown phase. The experimental configuration has two homodyne detections, (D_1, D_2) being one and (D_3, D_4) the other, such that the measurements are sensitive to phase changes in $u(\theta)$. The geometry of the setup is shown in Figure 1. In Figure 1 BS1, BS2 and BS3 are beam splitter mirrors that will reflect 50% of the incident photons and let 50% of them pass. When photons are reflected, the mirrors add a phase of $\pi/2$ to the expected density, while no phase is added to the expected density when photons pass through BS1, BS2 or BS3. It is easy to devise a way to have the expected density of photons changed by a $\pi/2$ phase by just delaying the photons that are reflected, and hence have interacted with the mirror, by a time $T/4$, where T is the period of the photon source. We will look for correlations between the pairs of photon detectors (D_1, D_2) and (D_3, D_4) .

The expected density of \pm -photons, generated by the source $u(\theta)$ is

$$h_{\pm}^u(\theta) = \frac{\beta}{2}(1 \pm \cos(\omega t + \theta)). \quad (10)$$

The expected density coming from $u(\theta)$ at each detector is

$$h_{\pm}^{D_1}(\theta) = \frac{\beta}{8} \left(1 \pm \cos \left(\omega t + \theta + \frac{\pi}{2} \right) \right), \quad (11)$$

$$h_{\pm}^{D_2}(\theta) = \frac{\beta}{8} (1 \pm \cos(\omega t + \theta + \pi)), \quad (12)$$

$$h_{\pm}^{D_3}(\theta) = \frac{\beta}{8} \left(1 \pm \cos \left(\omega t + \theta + \frac{\pi}{2} \right) \right), \quad (13)$$

$$h_{\pm}^{D_4}(\theta) = \frac{\beta}{8} (1 \pm \cos(\omega t + \theta)). \quad (14)$$

Note that we neglected factors of the form $\mathbf{k} \cdot \mathbf{x}$ coming from path contributions to the phase. We can do so considering the problem completely symmetric and remembering that only phase differences are relevant for the measurements we are proposing.

In similar fashion, the expected density of \pm -photons generated by the coherent sources α_i , with phase θ_i and amplitude $\alpha/2$, and α_j , with phase θ_j and amplitude $\alpha/2$, in each detector, is given by the following expressions:

$$h_{\pm}^{D_1}(\theta_i) = \frac{\alpha}{4} \left(1 \pm \cos \left(\omega t + \theta_i + \frac{\pi}{2} \right) \right), \quad (15)$$

$$h_{\pm}^{D_2}(\theta_i) = \frac{\alpha}{4} (1 \pm \cos(\omega t + \theta_i)), \quad (16)$$

$$h_{\pm}^{D_3}(\theta_j) = \frac{\alpha}{4} (1 \pm \cos(\omega t + \theta_j)), \quad (17)$$

$$h_{\pm}^{D_4}(\theta_j) = \frac{\alpha}{4} \left(1 \pm \cos \left(\omega t + \theta_j + \frac{\pi}{2} \right) \right), \quad (18)$$

where we again neglected path contributions to the phase and considered only the relevant phase at the detectors.

We should point out that in Eqs. (10)-(18), α and β are split in half at each semi-mirror, because each time a photon reaches a mirror there is a probability of 1/2 that the photon passes through and a probability of 1/2 that the photon is reflected by the mirror.

The probability of absorption in each detector, consistent with Eq. (7), is given most simply by the following equations. Some alternatives are formulated in Eqs. (39)-(46) at the end of this section.

$$p_{D_1} = \frac{C}{4} \left(2 + \cos \left(\omega t + \theta_i - \frac{\pi}{2} \right) + \cos \left(\omega t + \theta + \frac{\pi}{2} \right) \right), \quad (19)$$

$$p_{D_2} = \frac{C}{4} (2 + \cos(\omega t + \theta_i + \pi) + \cos(\omega t + \theta)), \quad (20)$$

$$p_{D_3} = \frac{C}{4} \left(2 + \cos(\omega t + \theta_j) + \cos \left(\omega t + \theta + \frac{\pi}{2} \right) \right), \quad (21)$$

$$p_{D_4} = \frac{C}{4} \left(2 + \cos \left(\omega t + \theta_j + \frac{\pi}{2} \right) + \cos(\omega t + \theta) \right), \quad (22)$$

where C is a constant that corresponds to the efficiency of the detection process.

The expected number of \pm photons in each detector is given, according to Eq. (8), by the following expressions:

$$E_t(D_1^\pm) = \left\langle \left(h_\pm^{D_1}(\theta_i) + h_\pm^{D_1}(\theta) \right) p_{D_1} \right\rangle, \quad (23)$$

$$E_t(D_2^\pm) = \left\langle \left(h_\pm^{D_2}(\theta_i) + h_\pm^{D_2}(\theta) \right) p_{D_2} \right\rangle, \quad (24)$$

$$E_t(D_3^\pm) = \left\langle \left(h_\pm^{D_3}(\theta_j) + h_\pm^{D_3}(\theta) \right) p_{D_3} \right\rangle, \quad (25)$$

$$E_t(D_4^\pm) = \left\langle \left(h_\pm^{D_4}(\theta_j) + h_\pm^{D_4}(\theta) \right) p_{D_4} \right\rangle. \quad (26)$$

Equations (23)-(26) use the fact that the expected number of photons at a detector is simply the sum of the number of photons from all sources. Also, in the equations above $\langle F(t) \rangle = \frac{1}{T} \int_0^T F(t) dt$, represents a time average of the random variable $F(t)$, where t is time and $(0, T)$ is a time interval such that $\omega T \gg 1$. It is straightforward to obtain the expressions for the total expected number of photons in each detector from Eqs. (3), (9), (15)-(18), (19)-(22), and (23)-(26), which we write as I_k , for $k = 1, \dots, 4$, with I_k a function of θ and θ_i or θ_j :

$$I_1 = |E_t(D_1^+) - E_t(D_1^-)| = \frac{C}{16} \left| \alpha + \frac{\beta}{2} + \left(\alpha + \frac{1}{2}\beta \right) \cos(\theta - \theta_i) \right|, \quad (27)$$

$$I_2 = |E_t(D_2^+) - E_t(D_2^-)| = \frac{C}{16} \left| \alpha + \frac{\beta}{2} - \left(\alpha + \frac{1}{2}\beta \right) \cos(\theta - \theta_i) \right|, \quad (28)$$

$$I_3 = |E_t(D_3^+) - E_t(D_3^-)| = \frac{C}{16} \left| \alpha + \frac{\beta}{2} - \left(\alpha + \frac{1}{2}\beta \right) \sin(\theta - \theta_j) \right|, \quad (29)$$

$$I_4 = |E_t(D_4^+) - E_t(D_4^-)| = \frac{C}{16} \left| \alpha + \frac{\beta}{2} + \left(\alpha + \frac{1}{2}\beta \right) \sin(\theta - \theta_j) \right|. \quad (30)$$

The expressions on the right hand side of (27)-(30) are nonnegative, independent of taking their absolute value, and so we subsequently drop the absolute values.

We are interested in the correlation between the two pairs of detectors. First we need the variances

$$\text{Var}_\theta(I_1 - I_2) = E_\theta((I_1 - I_2)^2) - (E_\theta(I_1 - I_2))^2, \quad (31)$$

$$\text{Var}_\theta(I_3 - I_4) = E_\theta((I_3 - I_4)^2) - (E_\theta(I_3 - I_4))^2, \quad (32)$$

and covariance

$$\text{Cov}_\theta((I_1 - I_2)(I_3 - I_4)) = E_\theta((I_1 - I_2)(I_3 - I_4)) - E_\theta(I_1 - I_2)E_\theta(I_3 - I_4), \quad (33)$$

where $E_\theta(I_k) = \frac{1}{2\pi} \int_0^{2\pi} I_k d\theta$, for $k = 1, \dots, 4$, is an expectation with respect to θ , with θ uniformly distributed on $[0, 2\pi]$. Thus

$$\text{Var}_\theta(I_1 - I_2) = \frac{1}{512} C^2 (\beta + 2\alpha)^2, \quad (34)$$

$$\text{Var}_\theta(I_4 - I_3) = \frac{1}{512} C^2 (\beta + 2\alpha)^2, \quad (35)$$

and

$$\text{Cov}_\theta((I_1 - I_2)(I_3 - I_4)) = -\frac{1}{512} C^2 (\beta + 2\alpha)^2 \sin(\theta_i - \theta_j). \quad (36)$$

The correlation is given by

$$\rho_\theta(I_1 - I_2, I_3 - I_4) = \frac{\text{Cov}_\theta((I_1 - I_2)(I_3 - I_4))}{\sqrt{\text{Var}_\theta(I_1 - I_2)\text{Var}_\theta(I_3 - I_4)}}, \quad (37)$$

or

$$\rho_\theta(I_1 - I_2, I_3 - I_4) = -\sin(\theta_i - \theta_j). \quad (38)$$

It is easy to show that (38) violates Bell's inequalities when four appropriate phases are chosen. Moreover, by making the source weak we can consider photon counts directly rather than intensities, and thus interpret the correlations as correlations of photon counts.

An attentive reader may object to our expression for the probability of detection, because we assume that the detector has the same probability to oscillate in phase with the noncoherent source as it has to oscillate with the coherent source, and that may bring some non-local characteristics to the model. We can respond to this by examining the following probability for absorption:

$$p_{D_1} = \frac{C}{2} \left(1 + \cos \left(\omega t + \theta_i + \frac{\pi}{2} \right) \right), \quad (39)$$

$$p_{D_2} = \frac{C}{2} (1 + \cos(\omega t + \theta_i)), \quad (40)$$

$$p_{D_3} = \frac{C}{2} (1 + \cos(\omega t + \theta_j)), \quad (41)$$

$$p_{D_4} = \frac{C}{2} \left(1 + \cos\left(\omega t + \theta_j + \frac{\pi}{2}\right) \right). \quad (42)$$

The probabilities above have no term in θ , but depend only on the phase of the coherent sources. This fact has the effect of wiping out all influences that the non-coherent source has on the detectors, and hence putting only local parameters, like θ_i or θ_j , depending on the detector, in the probability for detection. If we redo the computations for the correlation with the probabilities above, we end up with the same correlation function for a pair of homodyne detections. In fact, to point out the robustness of the result in face of the choice of probability for detection, we may examine the following set of probabilities.

$$p_{D_1} = \frac{C}{2} \left(1 + \frac{\alpha \cos(\omega t + \theta_i + \frac{\pi}{2}) + \beta \cos(\omega t + \theta + \frac{\pi}{2})}{\alpha + \beta} \right), \quad (43)$$

$$p_{D_2} = \frac{C}{2} \left(1 + \frac{\alpha \cos(\omega t + \theta_i + \pi) + \beta \cos(\omega t + \theta)}{\alpha + \beta} \right), \quad (44)$$

$$p_{D_3} = \frac{C}{2} \left(1 + \frac{\alpha \cos(\omega t + \theta_j) + \beta \cos(\omega t + \theta + \frac{\pi}{2})}{\alpha + \beta} \right), \quad (45)$$

$$p_{D_4} = \frac{C}{2} \left(1 + \frac{\alpha \cos(\omega t + \theta_j - \frac{\pi}{2}) + \beta \cos(\omega t + \theta)}{\alpha + \beta} \right). \quad (46)$$

The above expressions would have a different physical interpretation from the previous two presented. Each phase is given a probability that is proportional to the amplitude of the source with the corresponding phase. The stronger the source, the more probable to find the detector with the same phase. It is again easy to show that if we use these probabilities we get the same correlations as before. A detailed analysis of the selection of phase cannot be pursued here, but can be modeled by resonance processes dealing with the interaction of a photon and its absorbing atom.

3. CONCLUDING REMARKS

The significance of our results is this. Using a realistic local model of photons we derived the correlation $-\sin(\theta_i - \theta_j)$, identical

to that of quantum mechanics. That our model is local can be seen as follows (here we follow [1] and [5]). Locality requires

$$E(X|\theta_1, \theta_2, \theta) = E(X|\theta_1, \theta). \quad (47)$$

It is obvious that (47) follows immediately from subtracting (28) from (27), and observing the result does not depend on θ_2 , and similarly for the other cases. Equation (47) says simply that whatever is the result of the measurement at one homodyne detector, it must depend only on θ , the hidden variable, and the phase associated to this particular detector, and cannot be influenced by the phase at the other detector.¹

What our model shares with the quantum mechanical one given in Tan *et al.* [10] is the same phenomena of superposition of amplitudes, i.e., interference, as expressed by equations (23)-(26). By appropriate choice of the parameter β in equations (27)-(30) and integrating with respect to θ we obtain the same intensities I_i for the four detectors. Moreover, our local model is a relativistic one, because photons move with the speed of light and all interactions are local. However, the correlation they obtained was weaker than the ones we have. This is because they could not have maximum interference visibility, while our model makes a simplifying assumption that we can represent the mean distribution for a single photon by a coherent distribution. One expects that a more realistic model would also change the visibility, and give results closer to that obtained by Tan *et al.*

It is not uncommon for someone to ask of the kind of model considered here, "What happens when the source $u(\theta)$ is very weak and produces only one photon at a time?" Our response is that for the Tan *et al.* experiment the two coherent sources $\alpha(\theta_1)$ and $\alpha(\theta_2)$ provide a steady source of photons with which the occasional photon from $u(\theta)$ can locally interfere.²

ACKNOWLEDGMENTS

J. A. B. acknowledges support from the Department of Fields and Particles (DCP) and the Department of High Energy Physics

¹ This is a paraphrase of Bell's own words "It is the requirement of locality, or more precisely that the result of a measurement on one system be unaffected by operations on a distant system with which it has interacted in the past, that creates the essential difficulty. ... The vital assumption is that the result B for particle 2 does not depend on the setting a , of the magnet for particle 1, nor A on b ." [1]

² Other than the present paper, the authors are not aware of any other local realistic models that violate Bell's inequalities.

and Cosmology (Lafex) of the Brazilian Center for Physics Research (CBPF/CNPq). A. S. S. acknowledges financial support from CNPq (Brazilian Government's Support Agency).

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